

PRECISE ASYMPTOTIC OF EIGENVALUES OF RESONANT QUASILINEAR SYSTEMS

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ABSTRACT. In this work we study the sequence of variational eigenvalues of a system of resonant type involving p - and q -laplacians on $\Omega \subset \mathbb{R}^N$, with a coupling term depending on two parameters α and β satisfying $\alpha/p + \beta/q = 1$. We show that the order of growth of the k^{th} eigenvalue depends on $\alpha + \beta$, $\lambda_k = O(k^{\frac{\alpha+\beta}{N}})$.

1. INTRODUCTION

This paper is devoted to the study of the asymptotic behavior of eigenvalues of resonant quasilinear systems

$$(1.1) \quad \begin{cases} -\Delta_p u = \lambda \alpha |u|^{\alpha-2} u |v|^\beta \\ -\Delta_q v = \lambda \beta |u|^\alpha |v|^{\beta-2} v, \end{cases} \quad \text{in } \Omega$$

with Dirichlet boundary condition

$$(1.2) \quad u(x) = v(x) = 0 \quad \text{on } \partial\Omega.$$

Here, $\Omega \subset \mathbb{R}^N$ is a bounded domain with smooth boundary $\partial\Omega$, the s -Laplace operator is $\Delta_s u = \operatorname{div}(|\nabla u|^{s-2} \nabla u)$, the exponents satisfy $1 < p, q < +\infty$, and the positive parameters α, β satisfy

$$(1.3) \quad \frac{\alpha}{p} + \frac{\beta}{q} = 1.$$

The study of resonant systems has deserved a great deal of attention in recent years, and we may cite the works of Boccardo and de Figueiredo [5], Manasevich and Mawhin [24], Felmer, Manasevich and de Thèlin [16], Stavrakakis and Zographopoulos [26], among several others.

In several applications, such as bifurcation problems, anti-maximum principles, and existence or non-existence of solutions (see for example [4, 14, 16, 21, 26, 27, 28]) it is desirable to have precise bounds on the eigenvalues. In general this information is not well understood for elliptic systems, except for the first or principal eigenvalue. Several properties of this first eigenvalue were analyzed (existence, uniqueness, positivity, and isolation in bounded or unbounded domains, with different boundary conditions and with or without weights) and we refer the interested reader to [1, 12, 17, 28] among others.

Let us recall briefly that the existence of a sequence of variational eigenvalues for problem (1.1)–(1.2) was proved in [10], and the values λ_k are defined as

$$\lambda_k := \inf_{C \in \mathcal{C}_k} \sup_{(u,v) \in C} \frac{\frac{1}{p} \int_{\Omega} |\nabla u|^p dx + \frac{1}{q} \int_{\Omega} |\nabla v|^q dx}{\int_{\Omega} |u|^\alpha |v|^\beta dx},$$

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where \mathcal{C}_k is the class of compact symmetric ($C = -C$) subsets of $W_0^{1,p}(\Omega) \times W_0^{1,q}(\Omega)$ of (Krasnoselskii) genus greater or equal than k .

Throughout this work, the eigenvalues are counted repeated according to their multiplicity. We say that λ_k has multiplicity r if $\lambda_{k-1} < \lambda_k = \lambda_{k+1} = \dots = \lambda_{k+r-1} < \lambda_{k+r}$. In this case, it is a well know fact that the set of eigenfunctions corresponding to λ_k has genus greater or equal than r (see, for instance, [23]).

In the case of a single equation the existence of a sequence of variational eigenvalues together with the correct order of growth for these eigenvalues was first obtained by Garcia-Azorero and Peral in [23]. The constants in the asymptotic behavior were improved by Friedlander in [22]. Let us note that for the one-dimensional problem, these bounds can be refined, see [19, 20].

As for elliptic systems, the asymptotic growth of the eigenvalues is less understood even in the linear case. We may cite here the work of Protter [25], and also the works of Cantrell and Cosner [6, 7, 8] were lower bounds for the first eigenvalue were obtained. The exception for the lack of results in this direction comes from linear and nonlinear elasticity theory (see the survey of Antman [3]).

For nonlinear elliptic systems, up to our knowledge, the first work where this problem was addressed was [11] where the authors obtained a generalization of the Lyapunov inequality together with an upper bound of the variational eigenvalues in terms of the ones of a single p -Laplace equation for the one dimensional case. Later, in [18] we obtained lower and upper bounds for the eigenvalues of problem (1.1)–(1.2) in terms of the eigenvalues of a single p -Laplace and q -Laplace equations in any dimension $N \geq 1$. More precisely, for each k we prove that

$$c \left(\frac{k}{|\Omega|} \right)^{\frac{q}{N}} \leq \lambda_k \leq C \left(\frac{k}{|\Omega|} \right)^{\frac{p}{N}}$$

for suitable positive constants c, C depending on p and q .

We refer the interested reader to the introductions of [11, 18] for more information and references about the eigenvalues of quasilinear elliptic equations and resonant systems.

However, our previous bounds fail to reflect the coupling strength of the system which is given by the parameters α and β . Formally, by taking $\alpha = p$ and $\beta = 0$, we obtain a single p -Laplace equation replacing the system, namely

$$-\Delta_p u = \lambda p |u|^{p-2} u,$$

and similarly, for $\alpha = 0$ and $\beta = q$ we have

$$-\Delta_q v = \lambda q |v|^{q-2} v.$$

Hence, the order of growth of our upper (resp., lower) bound given in [18] is sharp for the case $\alpha = p$ (resp., $\beta = q$), since coincides with the true upper (resp., lower) order of growth of the eigenvalues (see [23]). On the other hand, both orders does not hold simultaneously even for those limit cases.

We can suspect that there exists a smooth transition for the order of growth of the eigenvalues between both limiting cases, and the main result of this work is to prove it. Namely, our main theorem is

Theorem 1.1. *Let $\{\lambda_k\}_{k \in \mathbb{N}}$ be the sequence of variational eigenvalues of problem (1.1)–(1.2). Then, there exist positive constants $c < C$ depending on p, q, α, β and N , such that*

$$c \left(\frac{k}{|\Omega|} \right)^{\frac{\alpha+\beta}{N}} \leq \lambda_k \leq C \left(\frac{k}{|\Omega|} \right)^{\frac{\alpha+\beta}{N}}$$

when $k \rightarrow \infty$.

Observe that if one consider linear operators, or even the same p -Laplace operator in both equations (i.e. $p = q$), the coupling parameters α and β are not reflected in the asymptotics of the eigenvalues since $\alpha + \beta = p$ in this case. We believe that this fact may be the reason why this phenomenon was not discovered earlier.

The proof of Theorem 1.1 follows directly from Weyl-type bounds for the spectral counting function $N(\lambda)$ which gives the number of eigenvalues less than a given value, that is

$$N(\lambda) = \#\{k : \lambda_k \leq \lambda\}.$$

Theorem 1.1 is equivalent to the following asymptotic bound for $N(\lambda)$:

$$C^{-1} \lambda^{\frac{N}{\alpha+\beta}} \leq N(\lambda) \leq c^{-1} \lambda^{\frac{N}{\alpha+\beta}}.$$

Up to our knowledge, this is the first case in the literature where the coupling parameters of an elliptic system appear explicitly modifying the power on the asymptotic order of growth of the eigenvalues.

1.1. Organization of the paper. The paper is organized as follows. In Section §2 we review some facts about the eigenvalue problem for the single s -Laplace equation in one space dimension

$$-(|u'|^{s-2}u')' = |u|^{s-2}u,$$

which can be found for example in [13], or [15].

In Section §3 we prove Theorem 1.1 by using a scaling argument as in [22]. The main drawback of this approach is the fact that the constant on the asymptotic expansion remains unknown since they depend on the first Dirichlet eigenvalue and the second Neumann eigenvalue of problem (1.1)–(1.2) when Ω is the unit cube Q_1 .

So, we are left with the problem of finding lower and upper bounds for those eigenvalues, which is a problem of independent interest. Thus, we consider the one-dimensional problem

$$(1.4) \quad \begin{cases} -(|u'(x)|^{p-2}u'(x))' &= \lambda\alpha|u|^{\alpha-2}u|v|^\beta, \\ -(|v'(x)|^{q-2}v'(x))' &= \lambda\beta|u|^\alpha|v|^{\beta-2}v, \end{cases}$$

on the interval (a, b) , and we will focus on the Dirichlet boundary condition.

In Section §4 we prove the following lower bound for the first Dirichlet eigenvalue $\lambda_1(p, q)$ of the one dimensional problem (1.4):

Theorem 1.2. *Let $\Lambda_1(s)$ be the first Dirichlet eigenvalue of*

$$-(|\varphi'|^{s-2}\varphi')' = \Lambda|\varphi|^{s-2}\varphi$$

on (a, b) . Then,

$$\left(\alpha^{\frac{\alpha}{p}}\beta^{\frac{\beta}{q}}\right)^{-1} \left(\frac{2}{\pi_{\alpha+\beta}}\right)^{\alpha+\beta} \frac{\Lambda_1(\alpha+\beta)}{\alpha+\beta-1} \leq \lambda_1(p, q).$$

The proof follows by using the Lyapunov inequality obtained in [11].

In [11] we obtained an upper bound of the first eigenvalue of the one dimensional problem (1.4) in terms of the first eigenvalue of the single p -Laplace equation. Moreover, upper bounds were obtained for all the variational eigenvalues, namely

$$(1.5) \quad \lambda_k \leq \frac{\Lambda_k(p)}{p} \left[1 + \left(\frac{p}{q}\right)^{q+1} (\Lambda_k(p))^{(q-p)/p}\right].$$

Here, $\Lambda_k(p)$ stands for the k^{th} eigenvalue of the p -Laplace with Dirichlet boundary conditions. Let us note that (1.5) holds for the one dimensional problem. In the N -dimensional case, (1.5) holds only for the first eigenvalue.

In this paper we improve this explicit upper bound for the first eigenvalue, and we use it to obtain asymptotic bounds for the k^{th} eigenvalue of a system in $\Omega \subset \mathbb{R}^N$ depending on the eigenvalues of the $(\alpha + \beta)$ -Laplace operator.

Section §5 is devoted to the proof of the following upper bound of the first Dirichlet eigenvalue of problem (1.4):

Theorem 1.3. *Assume $q < p$. Then, with the notation of Theorem 1.2, we have,*

$$\lambda_1(p, q) \leq \Lambda_1(\alpha + \beta) \left[\frac{1}{p} + \frac{1}{q} \left(\frac{\pi_{\alpha+\beta}}{(\alpha + \beta - 1) \int_0^{\pi_{\alpha+\beta}} \sin_{\alpha+\beta}^{\alpha+\beta}(t) dt} \right)^{1 - \frac{q}{\alpha+\beta}} \right].$$

Here $\sin_s(x)$ is defined implicitly as

$$x = \int_0^{\sin_s(x)} \frac{dt}{(1 - t^s)^{1/s}}$$

and π_s is given by

$$\pi_s = 2 \int_0^1 \frac{dt}{(1 - t^s)^{1/s}}.$$

See section §2 for more details and see also the paper [13].

For the N -dimensional problem we prove the following

Theorem 1.4. *Let $\lambda_1(p, q)$ be the first eigenvalue of (1.1) in the unit cube $Q_1 \subset \mathbb{R}^N$. Assume again that $q < p$. Then we have the explicit upper bound*

$$\lambda_1(p, q) \leq C \left(\frac{N(s-1)}{p} \pi_s^p + (s-1)^{q/s} \pi_s^q \left(\int_0^1 |\phi_s(t)|^s dt \right)^{-N(1-q/s)} \right),$$

where $s = \alpha + \beta$, $C = \max\{1; N^{(p-2)/2}; N^{(q-2)/2}\}$ and $\phi_s(t)$ is the first eigenfunction of the one dimensional s -Laplace operator with Dirichlet boundary conditions in the unit interval, that is

$$\phi_s(t) = \sin_s(\pi_s t).$$

Finally, in Section §6 we close this paper with some remarks on possible improvements on our result and some open questions.

2. SOME KNOWN FACTS

In this Section we recall some previous results which will be needed in the rest of the paper.

2.1. Variational setting. The variational characterization of eigenvalues follows from the abstract theory developed by H. Amann (see [2]). In [15] the authors showed the existence of infinitely many eigenpairs, i. e. $(u, v) \in W_0^{1,p}(\Omega) \times W_0^{1,q}(\Omega)$ and $\lambda \in \mathbb{R}$ such that

$$\int_{\Omega} |\nabla u|^{p-2} \nabla u \nabla \varphi + |\nabla v|^{q-2} \nabla v \nabla \psi dx = \lambda \int_{\Omega} (\alpha |u|^{\alpha-2} u \phi |v|^{\beta} + \beta |u|^{\alpha} |v|^{\beta-2} v \psi) dx$$

for any test-function pair $(\phi, \psi) \in W_0^{1,p}(\Omega) \times W_0^{1,q}(\Omega)$.

It is convenient to work with the variational characterization of the eigenvalues, defined through the Rayleigh quotient,

$$(2.1) \quad \lambda_k = \inf_{C \in \mathcal{C}_k} \sup_{(u,v) \in C} \frac{\frac{1}{p} \int_{\Omega} |\nabla u|^p dx + \frac{1}{q} \int_{\Omega} |\nabla v|^q dx}{\int_{\Omega} |u|^{\alpha} |v|^{\beta} dx},$$

where \mathcal{C}_k is the class of compact symmetric ($C = -C$) subsets of $W_0^{1,p}(\Omega) \times W_0^{1,q}(\Omega)$ of (Krasnoselskii) genus greater or equal that k .

2.2. One dimensional case. For the one dimensional s -Laplace operator in $\Omega = (0, L)$

$$(2.2) \quad -(|u'|^{s-2}u')' = \Lambda|u|^{s-2}u$$

with Dirichlet boundary conditions, we have

$$(2.3) \quad \Lambda_k(s) = \inf_{C \in \mathcal{C}_k} \sup_{u \in C} \frac{\int_a^b |u'|^s dx}{\int_a^b |u|^s dx},$$

with $u \in W_0^{1,s}(0, L)$.

Here, all the eigenvalues and eigenfunctions can be found explicitly (see [15]):

Theorem 2.1 (Del Pino, Drabek and Manasevich, [13]). *The eigenvalues $\Lambda_k(s)$ and eigenfunctions $u_{s,k}$ of equation (2.2) on the interval $[0, L]$ are given by*

$$\Lambda_k(s) = (s-1) \frac{\pi_s^s k^s}{L^s},$$

$$u_{s,k}(x) = \sin_s(\pi_s kx/L).$$

Remark 2.2. It was proved in [15] that they coincide with the variational eigenvalues given by equation (2.3). However, let us observe that the notation is different in both papers.

The function $\sin_s(x)$ is the solution of the initial value problem

$$-(|u'|^{s-2}u')' = (s-1)|u|^{s-2}u$$

$$u(0) = 0, \quad u'(0) = 1,$$

and is defined implicitly as

$$x = \int_0^{\sin_s(x)} \frac{dt}{(1-t^s)^{1/s}}.$$

Moreover, its first zero is π_s , given by

$$\pi_s = 2 \int_0^1 \frac{dt}{(1-t^s)^{1/s}}.$$

Let us note that both $\sin_s(x)$ and $\sin'_s(x)$ satisfy

$$|\sin_s(x)| \leq 1, \quad |\sin'_s(x)| \leq 1,$$

due to the Pythagorean like identity

$$(2.4) \quad |\sin_s(x)|^s + |\sin'_s(x)|^s = 1.$$

Finally, let us observe that the following integral is a constant depending only on s :

$$\int_0^{\pi_s} \sin_s^s(t) dt = K(s).$$

2.3. The Spectral Counting Function. Given the sequence $\{\lambda_k\}_{k \in \mathbb{N}}$, we introduce the spectral counting function $N(\lambda)$ defined as

$$N(\lambda) = \#\{k: \lambda_k \leq \lambda\}.$$

To avoid confusion, we will use $N_{sys}(\lambda)$ or $N_s(\lambda)$ to denote the eigenvalue counting functions of the system and the s -Laplace respectively. If necessary, we will write $N(\lambda, \Omega)$ to denote explicitly the set Ω where the eigenvalue problem is considered, and even $N^D(\lambda)$ or $N^N(\lambda)$ to indicate the Dirichlet and Neumann boundary conditions.

The main tool in order to obtain the asymptotic expansion of $N(\lambda)$ is the classical Dirichlet-Neumann bracketing introduced by Courant [9] in a version due to [19]:

Proposition 2.3 ([19], Theorem 2.1). *Let $U_1, U_2 \in \mathbb{R}^N$ be disjoint open sets such that $(\overline{U_1 \cup U_2})^{int} = U$ and $|U \setminus U_1 \cup U_2| = 0$. Then,*

$$\begin{aligned} N^D(\lambda, U_1) + N^D(\lambda, U_2) &= N^D(\lambda, U_1 \cup U_2) \\ &\leq N^D(\lambda, U) \\ &\leq N^N(\lambda, U) \\ &\leq N^N(\lambda, U_1 \cup U_2) \\ &= N^N(\lambda, U_1) + N^N(\lambda, U_2). \end{aligned}$$

3. ESTIMATES FOR THE SPECTRAL COUNTING FUNCTION

Let us begin with the following scaling argument. We denote by Q_1 the unit cube in \mathbb{R}^N and by Q_t the scaled cube of side length t .

Lemma 3.1. *Let $\lambda_1^1(p, q)$ (resp., $\lambda_1^t(p, q)$) be the first Dirichlet eigenvalue of problem 1.1 when $\Omega = Q_1$ (resp., $\Omega = Q_t$). Then,*

$$t^{\alpha+\beta} \lambda_1^t(p, q) = \lambda_1^1(p, q).$$

Proof. First, observe that given any pair $(u, v) \in W_0^{1,p}(Q_1) \times W_0^{1,q}(Q_1)$, if we define

$$u_t = tu\left(\frac{x}{t}\right), \quad v_t = tv\left(\frac{x}{t}\right),$$

then, $(u_t, v_t) \in W_0^{1,p}(Q_t) \times W_0^{1,q}(Q_t)$. Also, it is clear that $(u, v) \mapsto (u_t, v_t)$ is a bijection between $W_0^{1,p}(Q_1) \times W_0^{1,q}(Q_1)$ and $W_0^{1,p}(Q_t) \times W_0^{1,q}(Q_t)$.

Now, by simple change of variables, we get

$$\frac{\frac{1}{p} \int_{Q_t} |\nabla u_t|^p dx + \frac{1}{q} \int_{Q_t} |\nabla v_t|^q dx}{\int_{Q_t} |u_t|^\alpha |v_t|^\beta dx} = \frac{1}{t^{\alpha+\beta}} \frac{\frac{1}{p} \int_{Q_1} |\nabla u|^p dx + \frac{1}{q} \int_{Q_1} |\nabla v|^q dx}{\int_{Q_1} |u|^\alpha |v|^\beta dx}.$$

Finally, the variational characterization of the first eigenvalue gives the desired result, namely

$$\lambda_1^t(p, q) = \frac{\lambda_1^1(p, q)}{t^{\alpha+\beta}}.$$

This finishes the proof. \square

Remark 3.2. The argument used in the proof shows that if (u, v) is an eigenfunction associated to $\lambda_1^1(p, q)$ then (u_t, v_t) is an eigenfunction associated to $\lambda_1^t(p, q)$.

By applying the very same argument we obtain the following relation between the second Neumann eigenvalues (recall that the first Neumann eigenvalue is $\mu_1^t(p, q) = 0$).

Lemma 3.3. *Let $\mu_2^1(p, q)$ (resp., $\mu_2^t(p, q)$) be the second Neumann eigenvalue of problem (1.1) when $\Omega = Q_1$ (resp., $\Omega = Q_t$). Then,*

$$\mu_2^t(p, q) = \frac{\mu_2^1(p, q)}{t^{\alpha+\beta}}.$$

Proof. The proof follows exactly as the one in Lemma 3.1 and the variational characterization of the second Neumann eigenvalue, namely

$$(3.1) \quad \mu_2^t(p, q) = \inf_{C_t \in \mathcal{C}_2} \sup_{(u_t, v_t) \in C_t} \frac{\frac{1}{p} \int_{\Omega} |\nabla u_t|^p dx + \frac{1}{q} \int_{\Omega} |\nabla v_t|^q dx}{\int_{\Omega} |u_t|^{\alpha} |v_t|^{\beta} dx}.$$

The proof is now completed. \square

Now we are ready to prove Theorem 1.1.

Proof of Theorem 1.1. For any λ fixed, we take a lattice of cubes of side length $t \ll 1$ in \mathbb{R}^N with t depending on λ .

First we derive a lower bound for $N(\lambda)$ (equivalently, an upper bound for λ_k). From Lemma 3.1, $\lambda_1^t(p, q) = t^{-(\alpha+\beta)} \lambda_1^1(p, q)$, and taking

$$t = \left(\frac{\lambda}{\lambda_1^1(p, q)} \right)^{-\frac{1}{\alpha+\beta}},$$

we have

$$\lambda_1^t(p, q) = \lambda$$

Hence, since the first eigenvalue $\lambda_1^t(p, q)$ has multiplicity two (observe that (u, v) is an eigenfunction if and only if $(u, -v)$ is an eigenfunction), we have $N^D(\lambda, Q_t) = 2$.

By using the Dirichlet Neumann bracketing, a lower bound for $N(\lambda)$ is given by $2M$, where M is the number of cubes of the lattice contained in Ω . Since

$$t^N M \rightarrow |\Omega|$$

when $t \rightarrow 0$, we have

$$N(\lambda) \geq 2M \sim 2 \frac{|\Omega|}{t^N} = 2|\Omega| \left(\frac{\lambda}{\lambda_1^1(p, q)} \right)^{\frac{N}{\alpha+\beta}}.$$

The upper bound for $\lambda_k(p, q)$ follows since

$$k = N(\lambda_k(p, q)) \geq \frac{2|\Omega|}{\lambda_1^1(p, q)^{\frac{N}{\alpha+\beta}}} \lambda_k(p, q)^{\frac{N}{\alpha+\beta}}.$$

Let us find an upper bound for $N(\lambda)$. We use the bound for the second Neumann eigenvalue proved in Lemma 3.3, $\mu_2^t(p, q) = t^{-(\alpha+\beta)} \mu_2^1(p, q)$. Hence, taking

$$t = \left(\frac{c\lambda}{\mu_2^1(p, q)} \right)^{-\frac{1}{\alpha+\beta}},$$

for any $c > 1$, we have

$$\mu_2^t(p, q) = c\lambda > \lambda.$$

Therefore, for each cube we have $N^N(\lambda, Q_t) = 1$.

By using again the Dirichlet Neumann bracketing, an upper bound for $N(\lambda)$ is given by M , where M is the number of cubes covering Ω . As before, we have

$$N(\lambda) \leq M \sim \frac{|\Omega|}{t^N} = |\Omega| \left(\frac{c\lambda}{\mu_2^1(p, q)} \right)^{\frac{N}{\alpha+\beta}}.$$

The lower bound for $\lambda_k(p, q)$ follows since

$$k = N(\lambda_k(p, q)) \leq \frac{|\Omega|}{\mu_2^1(p, q)^{\frac{N}{\alpha+\beta}}} (c\lambda_k(p, q))^{\frac{N}{\alpha+\beta}}.$$

The Theorem is proved. \square

4. A LOWER BOUND FOR THE FIRST EIGENVALUE

In this section we restrict ourselves to the one dimensional case and prove Theorem 1.2. For that purpose, we use the following Lyapunov inequality for systems proved in [11]:

Theorem 4.1 ([11], Theorem 1.5). *Let us assume that there exists a positive solution of the system*

$$\begin{cases} -(|u'(x)|^{p-2}u'(x))' &= f(x)|u|^{\alpha-2}u|v|^\beta \\ -(|v'(x)|^{q-2}v'(x))' &= g(x)|u|^\alpha|v|^{\beta-2}v \end{cases}$$

on the interval (a, b) , with Dirichlet boundary conditions. Then, we have that:

$$(4.1) \quad 2^{\alpha+\beta} \leq (b-a)^{\frac{\alpha}{p'}+\frac{\beta}{q'}} \left(\int_a^b f(x) dx \right)^{\frac{\alpha}{p}} \left(\int_a^b g(x) dx \right)^{\frac{\beta}{q}}$$

This result gives Theorem 1.2 after replacing $f(x) = \alpha\lambda_1(p, q)$ and $g(x) = \beta\lambda_1(p, q)$. In fact, we have

$$(4.2) \quad 2^{\alpha+\beta} \leq \alpha^{\frac{\alpha}{p}} \beta^{\frac{\beta}{q}} (b-a)^{\frac{\alpha}{p'}+\frac{\beta}{q'}+\frac{\alpha}{p}+\frac{\beta}{q}} \lambda_1(p, q),$$

and let us note that

$$\frac{\alpha}{p'} + \frac{\beta}{q'} + \frac{\alpha}{p} + \frac{\beta}{q} = \alpha \left(\frac{1}{p} + \frac{1}{p'} \right) + \beta \left(\frac{1}{q} + \frac{1}{q'} \right) = \alpha + \beta.$$

So,

$$2^{\alpha+\beta} \leq \alpha^{\frac{\alpha}{p}} \beta^{\frac{\beta}{q}} (b-a)^{\alpha+\beta} \lambda_1(p, q).$$

The desired result follows from this inequality and the explicit formula for $\Lambda_1(\alpha + \beta)$ in Theorem 2.1, since:

$$\frac{2^{\alpha+\beta}}{(b-a)^{\alpha+\beta}} \frac{\alpha + \beta - 1}{\alpha + \beta - 1} \left(\frac{\pi_{\alpha+\beta}}{\pi_{\alpha+\beta}} \right)^{\alpha+\beta} = \frac{1}{\alpha + \beta - 1} \left(\frac{2}{\pi_{\alpha+\beta}} \right)^{\alpha+\beta} \Lambda_1(\alpha + \beta),$$

and therefore

$$\frac{1}{\alpha + \beta - 1} \left(\frac{2}{\pi_{\alpha+\beta}} \right)^{\alpha+\beta} \Lambda_1(\alpha + \beta) \leq \alpha^{\frac{\alpha}{p}} \beta^{\frac{\beta}{q}} \lambda_1(p, q).$$

This completes the proof of Theorem 1.2. \square

5. AN UPPER BOUND FOR THE FIRST EIGENVALUE

In this section, we first consider the one dimensional case and prove Theorem 1.3. Then, we extend the bound to the n -dimensional case by means of comparing the eigenvalue of the p -Laplacian with that of the pseudo p -Laplacian.

For the upper bound of the first eigenvalue, we need to improve the bound given in [11].

Given the variational characterization of the first eigenvalue,

$$\lambda_1(p, q) \leq \frac{\frac{1}{p} \int_0^L |u'|^p dx + \frac{1}{q} \int_0^L |v'|^q dx}{\int_0^L |u|^\alpha |v|^\beta dx}.$$

We denote $s = \alpha + \beta$ and choose $u = v = \varphi_1$, which is a multiple of the first Dirichlet eigenfunction of the single equation

$$-(|w'|^{s-2}w')' = \Lambda|w|^{s-2}w,$$

with Dirichlet boundary conditions, that is

$$\varphi_1(x) = \frac{L}{\pi_s} \sin_s(\pi_s x/L).$$

Due to the Pythagorean-like identity (2.4) we have

$$|\varphi_1'(x)| \leq 1.$$

So, taking $u = v = \varphi_1$ in the Rayleigh quotient we get

$$\lambda_1(p, q) \leq \frac{\frac{1}{p} \int_0^L |\varphi_1'|^p dx + \frac{1}{q} \int_0^L |\varphi_1'|^q dx}{\int_0^L |\varphi_1|^s dx}$$

Now, as $q < p$ we have that $s = \alpha + \beta < p$. Moreover, as $|\varphi_1'(x)| \leq 1$, the inequality

$$|\varphi_1'(x)|^p \leq |\varphi_1'(x)|^s$$

holds. On the other hand, by Hölder's inequality we obtain

$$\int_0^L |\varphi_1'|^q dx \leq \left(\int_0^L |\varphi_1'|^s dx \right)^{\frac{q}{s}} L^{1-\frac{q}{s}}$$

Collecting all these facts, we have

$$\begin{aligned} \lambda_1(p, q) &\leq \frac{\frac{1}{p} \int_0^L |\varphi_1'|^s dx}{\int_0^L |\varphi_1|^s dx} + \frac{1}{q} \frac{\left(\int_0^L |\varphi_1'|^s dx \right)^{\frac{q}{s}} L^{1-\frac{q}{s}}}{\int_0^L |\varphi_1|^s dx} \\ &\leq \frac{1}{p} \Lambda_1(s) + \frac{1}{q} \Lambda_1(s)^{\frac{q}{s}} \frac{L^{1-\frac{q}{s}}}{\left(\int_0^L |\varphi_1|^s dx \right)^{1-\frac{q}{s}}} \end{aligned}$$

Now we need an upper bound for

$$\left(\frac{L}{\int_0^L |\varphi_1|^s dx} \right)^{1-\frac{q}{s}}.$$

Indeed, we can compute the integral explicitly, obtaining

$$\begin{aligned} \int_0^L |\varphi_1|^s dx &= \left(\frac{L}{\pi_s} \right)^s \int_0^L \sin_s^s(\pi_s x/L) dx \\ &= \left(\frac{L}{\pi_s} \right)^{s+1} \int_0^{\pi_s} \sin_s^s(t) dt \\ &= \left(\frac{L}{\pi_s} \right)^{s+1} K(s). \end{aligned}$$

Then,

$$\lambda_1(p, q) \leq \frac{1}{p}\Lambda_1(s) + \frac{1}{q}\Lambda_1(s)^{\frac{q}{s}} \left(\frac{\pi_s^{s+1}}{L^s K(s)} \right)^{1-\frac{q}{s}}.$$

Now, by using the explicit formula for $\Lambda_1(s)$, we have

$$\lambda_1(p, q) \leq \Lambda_1(s) \left[\frac{1}{p} + \frac{1}{q} \left(\frac{\pi_s}{(s-1)K(s)} \right)^{1-\frac{q}{s}} \right].$$

The proof of Theorem 1.3 is finished. \square

Now we go back to the N -dimensional problem and try to find an explicit upper bound for the first eigenvalue $\lambda_1(p, q)$ in the unit cube Q_1 . That is, we prove Theorem 1.4.

In order to do this, we recall the pseudo p -Laplace operator

$$\hat{\Delta}_p u = \sum_{i=1}^N (|u_{x_i}|^{p-2} u_{x_i})_{x_i}$$

and consider the eigenvalue problem associated to $\hat{\Delta}_p$

$$\nu_1 = \inf_{(u,v) \in W_0^{1,p}(Q_1) \times W_0^{1,q}(Q_1)} \frac{\int_{Q_1} \frac{1}{p} |\nabla u|_p^p + \frac{1}{q} |\nabla v|_q^q dx}{\int_{Q_1} |u|^\alpha |v|^\beta dx}$$

where $|x|_s^s = \sum_{i=1}^N |x_i|^s$.

By the equivalence of norms

$$\begin{aligned} |x|_2 &\leq N^{(s-2)/2} |x|_s & \text{if } p > 2, \\ |x|_2 &\leq |x|_s & \text{if } 1 < p \leq 2, \end{aligned}$$

we have that $\lambda_1(p, q) \leq C\nu_1$, where $C = \max\{1; N^{(p-2)/2}; N^{(q-2)/2}\}$. So we need to bound ν_1 .

Recall that the first eigenfunction of

$$-\hat{\Delta}_s u = \nu_1 |u|^{s-2} u \quad \text{in } Q_1$$

with Dirichlet boundary conditions can be computed explicitly by separation of variables.

In fact, let $\phi_s(x) = \sin_s(\pi_s x)$ be the first eigenfunction of the one dimensional s -Laplace operator with Dirichlet boundary conditions in the interval $(0, 1)$, then

$$w_s(x) = \prod_{i=1}^N \phi_s(x_i)$$

is the first eigenfunction of the N -dimensional pseudo s -laplacian in the cube Q_1 .

Now, as in the one dimensional case, we use (w_s, w_s) , with $s = \alpha + \beta$ as a test function for ν_1 to get

$$\nu_1 \leq \frac{\frac{N}{p} \left(\int_0^1 |\phi_s(t)|^p dt \right)^{N-1} \int_0^1 |\phi_s'(t)|^p dt}{\left(\int_0^1 |\phi_s(t)|^s dt \right)^N} + \frac{\frac{N}{q} \left(\int_0^1 |\phi_s(t)|^q dt \right)^{N-1} \int_0^1 |\phi_s'(t)|^q dt}{\left(\int_0^1 |\phi_s(t)|^s dt \right)^N}.$$

By the Pythagorean-like identity

$$|\sin_s(t)|^s + |\sin_s'(t)|^s = 1$$

we conclude that $|\sin_s(t)|, |\sin_s'(t)| \leq 1$.

Also, as $q < s < p$ we get $|\phi_s(t)|^p \leq |\phi_s(t)|^s$ and $|\phi_s'(t)|^p \leq \pi_s^{p-s} |\phi_s'(t)|^s$

So

$$\begin{aligned}\nu_1 &\leq \frac{N}{p} \pi_s^{p-s} (s-1) \pi_s^s + \frac{\frac{N}{q} \left(\int_0^1 |\phi_s(t)|^q dt \right)^{N-1} \int_0^1 |\phi'_s(t)|^q dt}{\left(\int_0^1 |\phi_s(t)|^s dt \right)^N} \\ &= \frac{N(s-1)}{p} \pi_s^p + \frac{\frac{N}{q} \left(\int_0^1 |\phi_s(t)|^q dt \right)^{N-1} \int_0^1 |\phi'_s(t)|^q dt}{\left(\int_0^1 |\phi_s(t)|^s dt \right)^N}.\end{aligned}$$

To bound the other term, we use Hölder's inequality to get

$$\begin{aligned}\left(\int_0^1 |\phi_s(t)|^q dt \right)^{N-1} &\leq \left(\int_0^1 |\phi_s(t)|^s dt \right)^{(N-1)q/s}, \\ \int_0^1 |\phi'_s(t)|^q dt &\leq \left(\int_0^1 |\phi'_s(t)|^s dt \right)^{q/s} = \left((s-1) \pi_s^s \int_0^1 |\phi_s(t)|^s dt \right)^{q/s}.\end{aligned}$$

So

$$\nu_1 \leq \frac{N(s-1)}{p} \pi_s^p + (s-1)^{q/s} \pi_s^q \left(\int_0^1 |\phi_s(t)|^s dt \right)^{-N(1-q/s)}$$

which is an explicit bound for ν_1 . \square

6. AN EXPLICIT LOWER BOUND ON THE MAIN THEOREM

In Section §3 we obtained the following asymptotic bounds for the spectral counting function:

$$\frac{2|\Omega|}{\lambda_1^1(p, q)^{\frac{N}{\alpha+\beta}}} \lambda^{\frac{N}{\alpha+\beta}} \leq N(\lambda) \leq \frac{|\Omega|}{\mu_2^1(p, q)^{\frac{N}{\alpha+\beta}}} (c\lambda)^{\frac{N}{\alpha+\beta}}.$$

An explicit lower bound was given in the previous section by means of the explicit upper bound for $\lambda_1^1(p, q)$.

Although the previous formula holds for any constant $c > 1$, it is convenient to take $c = 1 + \lambda^{-1}$, since in this case we can rewrite the upper bound as

$$\frac{2|\Omega|}{\mu_2^1(p, q)^{\frac{N}{\alpha+\beta}}} \lambda^{\frac{N}{\alpha+\beta}} + O(1).$$

We conjecture that a stronger result holds, namely,

$$N(\lambda) = c(\Omega, \alpha, \beta) \lambda^{\frac{N}{\alpha+\beta}} + o(\lambda^{\frac{N}{\alpha+\beta}}).$$

Indeed, the proof follows immediately if it is true that

$$\lambda_1^1(p, q) = 2\mu_2^1(p, q).$$

For a single equation, the factor 2 does not enter since the first eigenvalue is simple. In this case, the equality $\lambda_1 = \mu_2$ holds for a single equation in dimension one. Up to our knowledge, it is not known for systems, nor in the case $N > 1$ even for a single p -Laplace equation.

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